



31 August 2015

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Subj: ONR Grant N00014-13-1-0032, "Modeling of Mid-Frequency Reverberation in Very Shallow Water"

Encl: (1) Final report for subject grant with accompanying SF 298.

Enclosed is the final report titled: "Modeling of Mid-Frequency Reverberation in Very Shallow Water: A Green's Function Approach and Application to TREX2013 Data Analysis" with the corresponding SF298 form for the subject grant.

This document constitute the Final Technical Report and deliverable for ONR Grant N00014-13-1-0032.

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Modeling of Mid-Frequency Reverberation in Very Shallow Water: A Green's Function Approach and Application to TREX2013 Data Analysis

FINAL REPORT FOR N00014-13-1-0032

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Acknowledgements

This project benefited greatly from ONR-funded experiments in 2013 and 2014 and the PIs of those efforts: Drs. Todd Hefner and DJ Tang. This work was supported by the US Office of Naval Research, Ocean Acoustics, project N00014-13-1-0032, Program Manager Ray Soukup.

Contents

Abstract	1
Keywords.....	1
Objectives.....	2
Project Background.....	3
Modeling Approach	5
A first-order Green's function approach	5
A modified higher-order approach	7
Numerical Examples	9
Environmental scenarios and input parameters	9
Algorithm for numerical simulations in steps	10
Application to TREX13 Data Analysis.....	17
TREX environmental scenario and input parameters	17
TREX-specified algorithm for numerical simulations	18
TREX Reverberation analysis: tentative model-data comparisons	18
Concluding Remarks and Implications for Future Research	25
Publications	26
References	27

Abstract

The long-term goals of this research are to better understand and accurately model low-to mid-frequency reverberation in shallow water environment. Specific goals are to develop a model of reverberation for conditions (1 – 10 kHz, ~20 m water depth, ~ 10 km range) corresponding to the ONR Target and Reverberation Experiment performed on the spring 2013 (TREX2013), develop a code and conduct computer simulations with environmental inputs typical for the chosen location, and apply this model to analysis of available TREX2013 data. This report presents a Green's function modeling approach that allows fast estimations of volume reverberation in complex shallow water environments. A simplified first-order version of the approach is considered to show how far-field scattering solutions obtained for free space can be incorporated into reverberation in complicated bounded, range-dependent, and stratified environments. A higher-order modification of this approach is considered as well, using a MFSB (Multiple Forward Single Backscatter) approximation. Application to analysis of TREX2013 reverberation and tentative model-data comparisons are presented.

Keywords

Underwater acoustic propagation and reverberation, Green's function in a shallow water waveguide, volume and roughness scattering, data inversion, TREX13 data analysis and model-data comparisons.

Objectives

The long-term goals of this research are to better understand and accurately model low- to mid-frequency propagation and reverberation in shallow water environment. This would result in development of a robust model (potentially a tool) for prediction and interpretation of reverberation in complex (range-dependent, potentially 3D-dependent) waveguides. This report presents a Green's function modeling approach that allows fast estimations of propagation and reverberation in shallow water. The approach also addresses a basic science question: how scattering solutions obtained for free space (or measured in direct-path, short-range conditions) can be incorporated into long-range reverberation in bounded and complicated environments, such as shallow water or deep-water waveguides, where effects of multiple interactions with boundaries and volume heterogeneities must be taken into account, providing both sufficient accuracy and speed of calculations.

Specific goals of this research are to develop a model of reverberation for conditions (1 – 10 kHz, ~ 20 m water depth, ~ 10 km range) corresponding to the ONR Target and Reverberation Experiment performed on the spring 2013 (TREX2013), develop a code and conduct computer simulations with environmental inputs typical for the chosen location, and apply this model to analysis of available TREX2013 data, see [TREX13 web-site, TREX13 Workshop (2012, 2013, 2014)]. This report demonstrates capabilities of the developed codes based on Green's function modelling approach to provide extremely fast predictions of reverberation in a very shallow water environment, gives numerical examples for several relevant scenarios, discusses possibilities for a quantitative interpretation of reverberation data, presents tentative TREX2013 data-model comparisons, and suggests implications for future research.

Project Background

A proposal for this project was written in response to the ONR call for suggestions and comments on White Paper (2011) “Mid-frequency reverberation measurements with full companion environmental support”, which defined conditions for approaching reverberation experiment in a very shallow water (~20 m depth). Scattering mechanisms to be addressed were specified as well: sea surface and bottom roughness, heterogeneity in water column (near surface bubbles, fish schools) and sediments (e.g. large shells and mud inclusions). White Paper emphasized, however, importance to address applicability of this study to more general environments. In addition to the general problem of reverberation, specific basic science issues to be addressed were defined: (1) Mid-frequency sound scintillation index and coherent field, (2) Evolution of mid-frequency spatial coherence as a function of range from source, (3) Sediment sound speed and attenuation as a function of frequency, (4) Effects of spatial variability of bottom scattering strength, (5) Impact of sediment ripple fields on waveguide propagation and sediment volume scattering.

The proposal for this project, “Modeling of mid-frequency reverberation in very shallow water”, submitted in 2012, can be briefly described as follows. The main goal of the project was to enhance the modeling component of the approaching experiment, and to contribute to development of non-traditional approaches to shallow water reverberation, i.e. to certainly support the idea expressed in White Paper (2011), that a new area of investigation for WPRM (Wave Propagation in Random Media) in shallow water should be suggested. The proposal mentioned, however, that for such very shallow water and to support this specific experiment, the focus of such investigation should be broader than or even different from the one in the mentioned White Paper, where effects of only internal waves were emphasized in this connection. This project suggested that the investigation most relevant to very shallow water environment, where sediment is a critical part of the propagation channel, should be focused on the effect of 3D variability of the sediment properties along with corresponding effects of propagation and 3D refraction within the seafloor.

It was assumed that in such complicated conditions some multiple scattering effects may be important, and suggested to address these effects exploiting an approach developed by De Wolf (1971) and then successfully used by Ishimaru, Tatarskii, and others for description of electromagnetic propagation and multiple scattering in a turbulent atmosphere. Particularly, the approach is applicable to description of the so-called backscattering enhancement effect, known also as “weak localization” and “coherent backscatter” effects. This approach is well recognized in WPRM community, and called the MFSB-approximation (Multiple-Forward-Single-Backward-scatter). While similar terminology appears frequently in the underwater acoustics literature, such approach for conditions of shallow water reverberation has not yet been adequately developed, although experimental observations exist, Sabra (2010).

For description of multiple-scattering effects in acoustics of marine sediments, the MFSB-approach was used by Ivakin (1999), and then it was reformulated for the case of

reverberation in a shallow-water waveguide in Ivakin (2008). The main complication in this case (in comparison with previous work) appears because of the multi-path conditions for propagation in a waveguide. Recent developments suggest treating this complication by direct considering both the heterogeneous sediment and water column as parts of ocean waveguide Ivakin (2011, 2012). In this research, the approach was proposed to be developed further for a very shallow water environment. In particular, it was suggested to perform a numerical modeling for several typical scenarios including both water column and seabed as possible mechanisms of reverberation to better understand their relative effects.

One of practically important directions of this research is development, using both analytical and numerical modeling, of practical relationships between reverberation and environmental parameters, with a specific goal to evaluate the sensitivity of measured reverberation intensity to these parameters, to determine which environmental parameters or mechanisms of scattering are critical for reverberation in chosen acoustic and environmental very shallow-water scenarios. This may help to reduce the number of parameters needed to be measured or taken into account to provide “full environmental support” of the field experiment proposed in White Paper (2011). Another practically important application of this modeling would be evaluation of the possibility for inferring environmental parameters from reverberation data. New techniques and algorithms for such inferring were expected to be developed.

This project was proposed for a two-year period, FY2013 and FY2014. During the first- year, a simplified version of the reverberation model was planned to be developed, with a focus on bottom scattering mechanisms, to facilitate developing codes and conducting pre-test computer simulations, helping with planning the acoustic experiment and environmental ground truth measurements. Results of this work were presented in Ivakin (2013), and Hefner et al (2013a, 2013b). During the second year of work, in FY2014, a more general model of reverberation was planned to be developed and applied to analysis of TREX2013 data based on available by then acoustic and environmental inputs. Also, it was suggested to consider analytical expressions for reverberation, methods of its inversion for environmental parameters, and applicability of results obtained in specific TREX2013 conditions to more general shallow water scenarios. Some results of this work are presented in Ivakin (2015a), and Hefner et al (2015a, 2015b).

Modeling Approach

This report presents a Green's function modelling approach that allows fast estimations of reverberation in shallow water. The approach also addresses a basic science question: how far-field scattering solutions obtained for free space can be incorporated into reverberation in bounded and complicated environments, such as shallow water or deep-water waveguides, where effects of multiple scattering and interaction with boundaries and volume heterogeneities must be taken into account.

A basic approach, its simplified first-order version and initial results, are based on models discussed in Ivakin (2010, 2011, 2012, 2013). A higher-order modification of this approach exploits ideas of a MFSB (Multiple Forward Single Backscatter) approximation developed by De Wolf (1971) for studying electromagnetic propagation and multiple scattering effects in a turbulent atmosphere. For description of such effects in acoustics of marine sediments, the MFSB-approach was applied in Ivakin (1999), and reformulated for the case of reverberation in a shallow-water waveguide in Ivakin (2008). Its further development is presented in Ivakin (2015a).

A first-order Green's function approach

Consider the total field (acoustic pressure) as a sum of the unperturbed field and a first-order (single-scattered) field expressed through correspondent components of the full-field Green's function,

$$p = p_0 + p_1 = AG = A(G_0 + G_1) \quad (1)$$

with a coefficient, A , introduced so that Green's function near the source represents a spherical wave of the unite magnitude, i.e. $|\vec{r} - \vec{r}''|G(\vec{r}, \vec{r}'') \xrightarrow{\vec{r} \rightarrow \vec{r}''} 1$. Analytical results and numerical solutions for zeroth-order Green's function are currently available for rather complicated background (or unperturbed) media, such as range-dependent shallow water or deep-water waveguides, and are based on various propagation models, such as PE, normal modes, wave-number integration, or various ray-based approximations.

A solution for the single-scattered field is known, see e.g. Ivakin (2008, 2010, 2011, 2012, 2013), which then results in an integral expression, Ivakin (2011, 2012)

$$\begin{aligned} \langle I_1(\vec{r}', \vec{r}'') \rangle &= |A|^2 \int (\rho_0(\vec{r}') / \rho_0(\vec{r}))^2 M_V(\vec{r}) |G_0(\vec{r}, \vec{r}') G_0(\vec{r}, \vec{r}'')|^2 d^3r = \\ &= |A|^{-2} \int M_{Vef}(\vec{r}) I_0(\vec{r}, \vec{r}') I_0(\vec{r}, \vec{r}'') d^3r \end{aligned} \quad (2)$$

where $I_{0,1} = |p_{0,1}|^2$ are the zeroth- (unperturbed) and first-order (scattered) intensities, $\rho_0(\vec{r})$ is the density in unperturbed medium, $M_{Vef} = (\rho_0(\vec{r}') / \rho_0(\vec{r}))^2 M_V$, and M_V is a

volume scattering coefficient, or the scattering cross-section per unit volume. Limitations of Eq.(2) result from assumption that the spectrum of heterogeneity is a smooth enough function, allowing neglecting bistatic scattering effects by ignoring difference in wave vectors appearing in the stratified environment with both down-going and up-going waves, Ivakin (2011, 2012). Despite of its simplicity, Eq.(2) is rather general and can be applied to scattering near marine boundaries, seabed or sea surface, or near an arbitrary reference interface within marine environment, using a general procedure described in Ivakin (2012).

A specific expression for M_V can be obtained from a far-field solution for scattering in free space with a specified type of heterogeneity, continuum or discrete. For backscatter from continuum heterogeneity caused by local spatial fluctuations of the density and sound speed, we have, see Ivakin (2011, 2012),

$$M_V = 2\pi k^4 \Phi_{2\vec{k}}(\varepsilon_{\rho c}) \quad (3)$$

where $\varepsilon_{\rho c} = (\rho c)/(\rho_0 c_0) - 1$ are relative fluctuations of the impedance, and Φ is their power spectrum at the backscattering Bragg's wave-vector. In the case of discrete scatterers, such as particles and objects of various kinds, a general expression for the incoherent volume scattering strength of randomly distributed objects is of the form, see Ivakin(2012),

$$M_V = C_V \langle \sigma / v \rangle \quad (4)$$

where σ and v are the individual scattering cross section of the object and its volume, their ratio is averaged (given a fixed volume) over other parameters, such as size, shape and orientation, and $C_V \ll 1$ is the total volume concentration of the objects.

Therefore, the approach provides a simple way for calculations of reverberation as follows. An integral expression for the backscatter intensity, given by Eq.(2), has a factorized integrand comprised of two kernels, the two-way propagator and the scattering kernel. The propagator is a product of two local intensities, each being defined along one of the two ways of propagation. For rather complicated environments, generally depth-range-dependent, the local intensity can be calculated using available models and codes, e.g. PE, normal modes, ray-based, or other approximations.

The scattering kernel is defined as a local volume scattering coefficient and exploits simple first-order solutions for far-field scattering from a heterogeneous volume in free (unbounded) space. It can be specified for sea-water column and seabed with continuum heterogeneity, such as spatial fluctuations of density and sound speed, and/or discrete randomly and sparsely distributed targets, such as gas bubbles, fish, shells, lenses-like inclusions, oil droplets, solid hydrate particles, and others, Ivakin (2011, 2012).

Moreover, the scattering kernel can include a component due to contribution of roughness at an arbitrary number of interfaces at depths $z = z_j$ as follows

$$M_{Vef}(\vec{r}) = m^{-2} M_V + \sum_j M_{Rj} \delta(z - z_j) \quad (5)$$

$$M_{Rj} = \frac{k_o^4}{4} D_j^2 \Phi_{2\vec{k}}^{(R)}(\zeta_j) \quad (6)$$

$$D_j = \left(\frac{n_j}{m_j} - \frac{n_{j+1}}{m_{j+1}} \right) (n_j + n_{j+1}) \quad (7)$$

where $m = \rho / \rho_0$ and $n = c_0 / c$ are relative density and refraction index of the background medium, M_{Rj} , $\Phi^{(R)}$, and D_j are the roughness scattering coefficient, the power spectrum, and a contrast factor at j -th interface, respectively. In more detail, this case is described in Ivakin (1998, 2015c).

Note, that this approach provides estimations of bottom reverberation without calculations of the equivalent surface scattering strength (although may include it as a particular case, Ivakin (2012, 2013, 2015a). The approach is applicable, generally, for an arbitrary distribution of scatterers defined by their volume scattering coefficient M_v , and roughness scattering coefficient M_R , see Eqs.(3-5). Therefore, it allows a fast estimation of potential contributions of different mechanisms of scattering with arbitrary strengths and locations.

A modified higher-order approach

Integral expression (2) for the reverberation intensity can be generalized by treating the propagating kernel stochastically and accounting for multiple forward-scatter effects, using MFSB approximation developed by De Wolf (1971). According to this approach, the total field is presented as a sum of a multiple forward-scattered field and a single-backscattered field, or through correspondent components of the full Green's function,

$$p = p_f + p_s = AG = A(G_f + G_s) \quad (8)$$

Then one obtains the integral expression

$$\begin{aligned} \langle I_s(\vec{r}_1, \vec{r}_2) \rangle &= |A|^2 \int \left\langle M_{vef}(\vec{r}) \left| G_f(\vec{r}, \vec{r}_1) G_f(\vec{r}, \vec{r}_2) \right|^2 \right\rangle d^3r = \\ &= |A|^{-2} \int \left\langle M_{vef}(\vec{r}) I_f(\vec{r}, \vec{r}_1) I_f(\vec{r}, \vec{r}_2) \right\rangle d^3r \end{aligned} \quad (9)$$

where $I_{f,s} = |p_{f,s}|^2$ are the multiple forward- and single back-scattered intensities.

Despite of apparent similarity to the first-order expressions (1,2), Eqs.(8,9) represent a substantial modification and improvement. The modified approach allows taking into account such multiple scattering effects, as the scintillations of propagated intensity, the two-way propagation coherence, and the related “backscattering enhancement” known also as “coherent backscatter” and “weak localization”.

To show how and where this improvement appears and can be used, recall an important energy conserving advantage of the forward scattering propagator, resulting in $\langle I_f \rangle \approx I_0$. Then for the backscatter from uniform scatterer distributions one obtains

$$\langle I_b \rangle = \langle I_s(\vec{r}_1, \vec{r}_1) \rangle \approx |A|^{-2} \int M_{vef}(\vec{r}) I_0^2(\vec{r}, \vec{r}_1) [1 + K_I(\vec{r}, \vec{r}_1)] d^3r \quad (10)$$

where K_I is the scintillation index. The effect of backscattering enhancement appears here automatically because $K_I > 0$.

Consider now applicability of this modified approach to analysis of the seabed reverberation in complicated conditions, such as in the TREX13 environment, where marine sediments have a substantial lateral variability. To apply Eq.(9), a general procedure can be used similar to described in Ivakin (2012). Consider spatial variables of integration in (9) of the form $\vec{r} = (\vec{r}_\perp, \xi)$, $\vec{r}_\perp = \vec{r}|_{\xi=0}$, and introduce normalizing functions, for example as incident forward propagation wave functions on the bottom surface, $G_{0f}(\vec{r}_\perp, \vec{r}_2)$, so that $G_f(\vec{r}, \vec{r}_{1,2}) = G_{0f}(\vec{r}_\perp, \vec{r}_{1,2}) g_{1,2}(\vec{r})$. Then Eq.(9) can be presented in the form

$$\langle I_s(\vec{r}_1, \vec{r}_2) \rangle = |A|^2 \int \langle M_s(\vec{r}_\perp) \rangle \left| \left\langle G_{0f}(\vec{r}_\perp, \vec{r}_1) G_{0f}(\vec{r}_\perp, \vec{r}_2) \right| \right|^2 d^3r \quad (11)$$

where M_R is a local equivalent interface scattering cross section, or dimensionless bottom scattering strength,

$$M_R = \int (\rho_1 / \rho)^2 |g_1 g_2|^2 M_v d\xi \quad (12)$$

Eqs.(9-12) can be used for analysis of volume reverberation in a complex TREX13 environment taking into account heterogeneity of both water column and the seabed. For instance, TREX13 sediments were very heterogeneous, and had a substantial lateral variability of the sediment type and composition, mostly due to variations in mud and shell content, Hefner and Tang (2014).

The observed variability of sediment composition (e.g. change from well-sorted medium sand to muddy or shelly sand) may result in noticeable variations in sound speed and attenuation. Particularly important in this case are changes in local critical angle, affecting local reflection, penetration, and scattering properties. For instance, alternations of the sediment type along the propagation paths may have a strong accumulation effect on the propagator kernel in (9-11), resulting in the propagation intensity scintillations and corresponding backscattering enhancement, see Eq.(10). The alternations across the path result in an effective averaging of the scattering kernel in (11) given by the local bottom scattering strength M_R , which is very sensitive to changes of the sediment type, especially at shallow grazing angles (below the critical angle). Therefore, such lateral variability may significantly affect the reverberation at long ranges. A quantitative analysis of the effect based on this approach is being planned as a part of future work.

Numerical Examples

Examples here are to demonstrate capabilities of the first-order Green's function modeling approach to provide fast estimations of reverberation, to show three critical steps in the algorithm of these estimations, and to help analyze potential contributions of different mechanisms of scattering in shallow water environments (heterogeneous water column, rough and heterogeneous bottom). Several scenarios with different types of a layered mud/sand bottom are considered, which are applicable to very shallow water, as well as to more general shallow water environments, see e.g. ONR Workshop (2015). Input parameters for bottom are chosen to be representative and typical for mud and sand sediments.

Environmental scenarios and input parameters

In the following numerical examples, Green's function magnitude, as a function of range and depth (specified for given very shallow water range-independent environment), was calculated using a PE code, Tang (2014), with input parameters as follows:

Frequency – 3.5 kHz;
Water depth – 20 m;
Source/receiver depth – 19 m;

Three bottom scenarios were chosen to represent sand/mud sediments with different thickness of the mud layer (h) placed above sand halfspace:

- (1) sand halfspace ($h=0$),
- (2) 0.5 m of mud above sand halfspace,
- (3) 3.5 m of mud above sand halfspace.

Calculations for the penetration field within the bottom were made for depths down to 5 m below bottom surface.

Acoustic parameters for water, mud, and sand are taken as follows:

	water	mud	sand
Sound speed [m/s]	1530	1500	1670
Density [g/cm^3]	1.0	1.6	2.0
Attenuation (loss tangent), δ	0	0.002	0.01

Numerical simulations in steps

The algorithm of fast estimations of shallow water reverberation includes three critical steps. The first step, according to Eq.(2), is to evaluate the propagation kernel as a function of locations of the source/receiver and scattering point. For this, the Green's function magnitude should be pre-calculated in the waveguide of interest. Importantly, the algorithm requires calculations only at one (central) frequency. These calculations are illustrated in Figures 1-3.

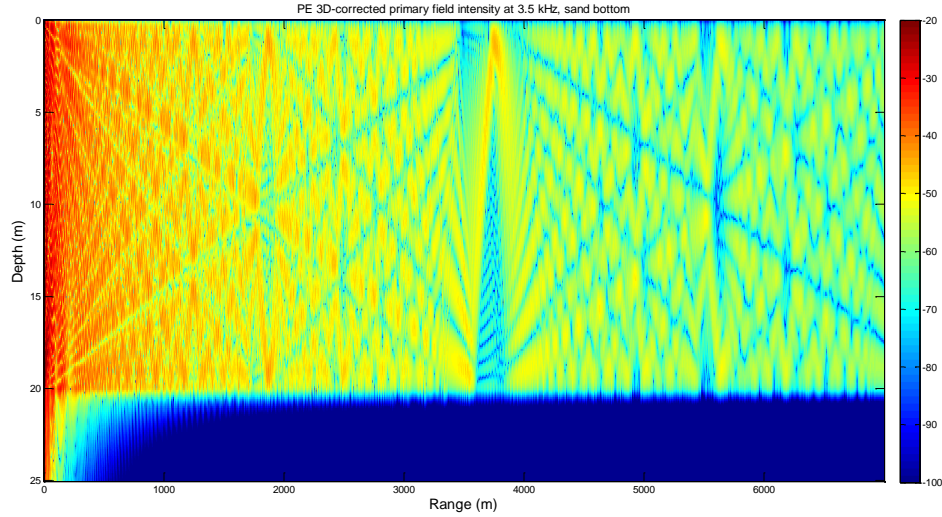


Figure 1: *Intensity of 3.5 kHz PE-propagation field in shallow water (of 20m depths) with sand bottom. The source is 1 m above bottom.*

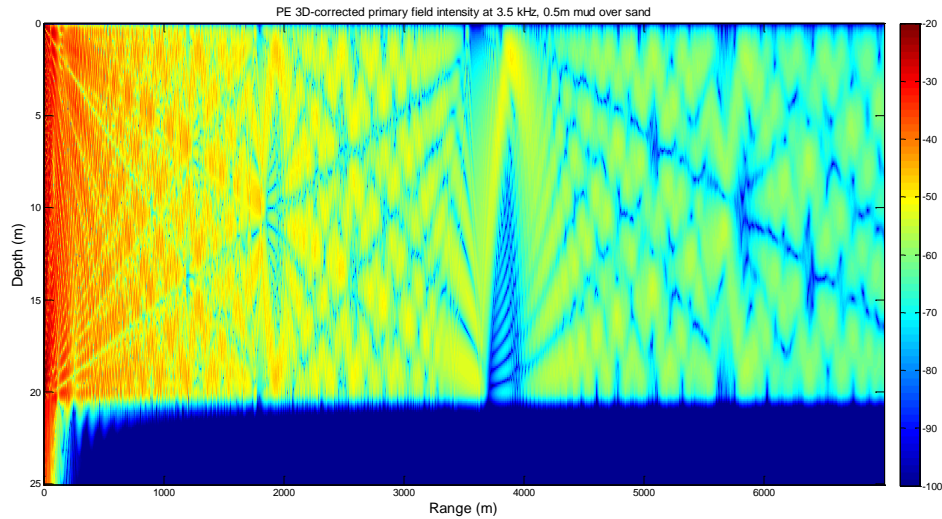


Figure 2: *Intensity of 3.5 kHz PE-propagation field in shallow water (20m depth) with a two-layer, 0.5m of mud over sand, bottom. The source is 1 m above bottom.*

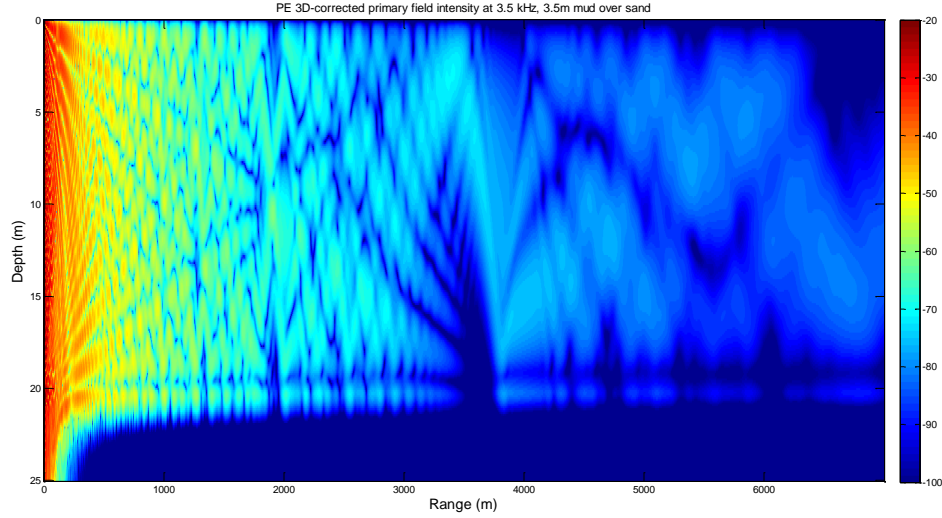


Figure 3: *Intensity of 3.5 kHz PE-propagation field in shallow water (20m depth) with a two-layer, 3.5m of mud over sand, bottom. The source is 1 m above bottom.*

The second step of the algorithm is illustrated in Figures 4-6. It is seen from Figure 4, that the range-dependence of the propagation intensity on the bottom surface, which is important for estimation of bottom reverberation, is a strongly oscillating function. For further analysis, such calculations require a kind of smoothing. For a particular kind, results are shown in Figure 5, where the smoothing was performed for a fixed range span $L = 100$ m. A simple rationale for such smoothing may result from possible uncertainty in ranges/depths, which can be substantial in some measurement techniques, e.g. in the case where source or/and receiver locations are not fixed. For instance, such technique was used in the TREX propagation measurements with a fixed vertical array in mid-water (at 9-15 m depths) and a CW-source moving (i.e. not fixed) at a 10 m depth, so that a range dependence of the Transmission Loss (TL) was measured. It was noticed that, for regions with sand bottom covered by a mud layer (about 20cm thick) and regions with sand bottom (free of mud), TL is essentially the same. This effect is easy to see in Figure 6, where even a 50 cm mud layer is shown to be not sufficient for a noticeable effect on propagation intensity averaged over 100 m ranges and 9-15 m depths spans. Another reason for smoothing over a range span will be discussed later in connection with TREX reverberation data analysis.

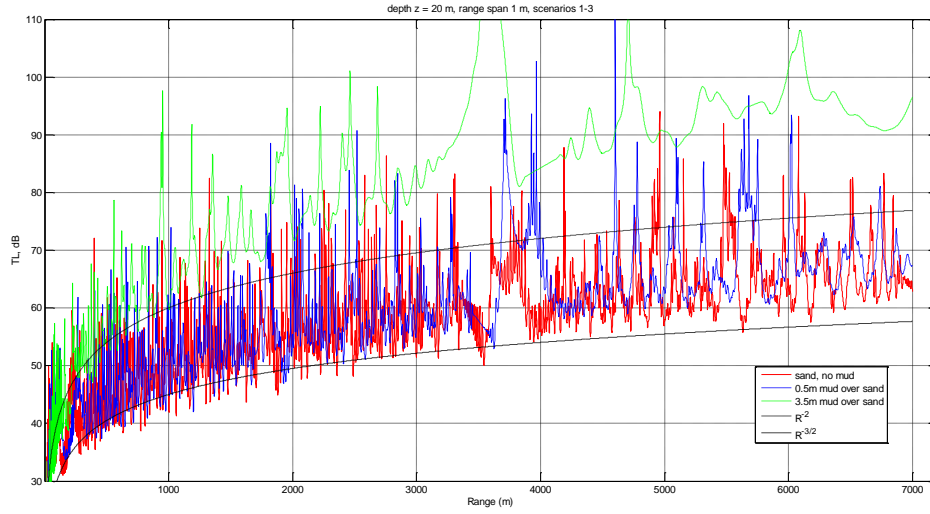


Figure 4: Range-dependence of Transmission Loss at the bottom surface (at 20m depth), 3.5 kHz PE-propagation field for three different types of bottom: sand half-space (red), 0.5m of mud over sand (blue), and 3.5m of mud over sand half-space (green). The source is 1 m above bottom. Also shown are $R^{3/2}$ and R^2 range-dependences.

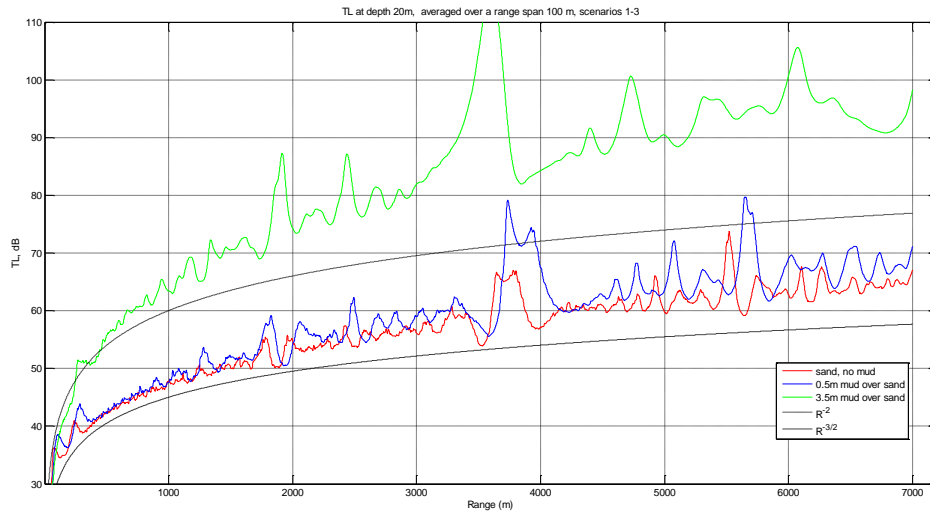


Figure 5: Smoothed over 100m range span range-dependence of Transmission Loss at the bottom surface (at 20m depth), 3.5 kHz PE-propagation field for three different types of bottom: sand half-space (red), 0.5m of mud over sand (blue), and 3.5m of mud over sand half-space (green). The source is 1 m above bottom. Also shown are $R^{3/2}$ and R^2 range-dependences.

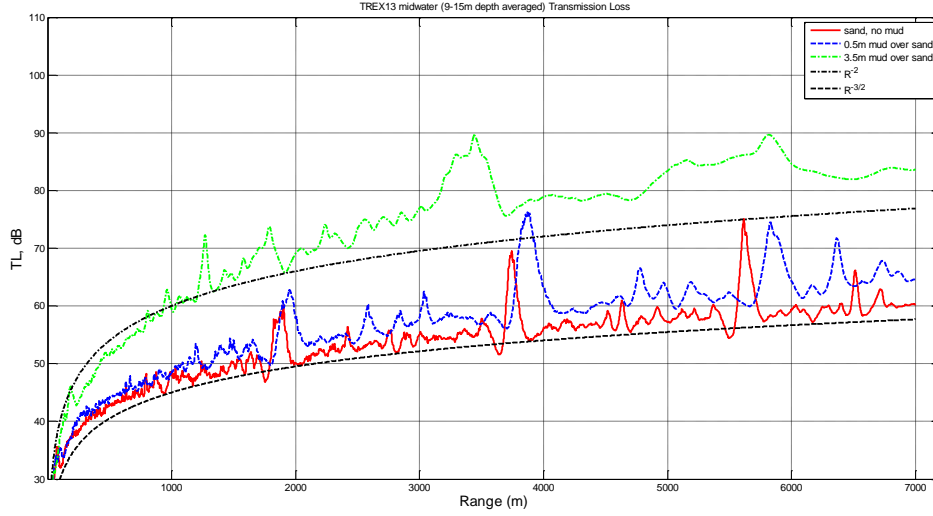


Figure 6: Range-dependence of Transmission Loss in mid-water, averaged (smoothed) over 100m range span and over 9-15m depths, 3.5 kHz PE-propagation field in shallow water (20m depth) for three bottom types: sand half-space (red), 0.5m of mud over sand (blue), and 3.5m of mud over sand half-space (green). The source is 1 m above bottom. Also shown are $R^{3/2}$ and R^2 range-dependences.

The third step of the algorithm provides extremely fast evaluation of reverberation in a shallow water waveguide using pre-calculated for this waveguide range-depth dependence of single-frequency Green's function intensity. The effect of the frequency bandwidth is considered at this step as follows. According to Eq.(2), the integration should be performed over a horizontal cross-section of the scattering (ensonified) volume, which assumes a span of azimuthal angles, $\Delta\varphi$, (to account for possible azimuthal directivity, e.g. for horizontal arrays, such as FORA). Also, the integration over scattering area presumes another span of ranges, Δr , which may be significantly different from above mentioned smoothing scale L . The span Δr is related to a span of double (two-way) propagation time, and can be defined (although only approximately) in the same way as in a direct-path (short-range) scenario, through the frequency bandwidth, Δf , or through the pulse duration, τ , as follows: $\Delta r = c\tau/2 = c/\Delta f/2$. In the following calculations, input parameters needed at this step are as follows: $\Delta f = 100$ Hz, $\Delta\varphi = 2.6^\circ$. This gives an estimate $\Delta r = 8$ m.

At this step, reverberation caused by bottom volume heterogeneity and rough interfaces, as well as volume reverberation from scatterers in water column, can be calculated, generally, for an arbitrary distribution of scatterers defined by their volume scattering coefficient M_v , and roughness scattering coefficient M_R , see Eqs.(3-7). Therefore, this allows fast estimating and comparing relative contributions of different scattering mechanisms, with different nature, strengths, and locations.

Results for chosen bottom scenarios are presented in Figures 7-11, showing reverberation caused by different types of scatterers: heterogeneities of water column (Figure 7), bottom volume (Figures 8, 9), and bottom interface roughness (Figures 10, 11), using representative values for input parametrs, M_V and M_R .

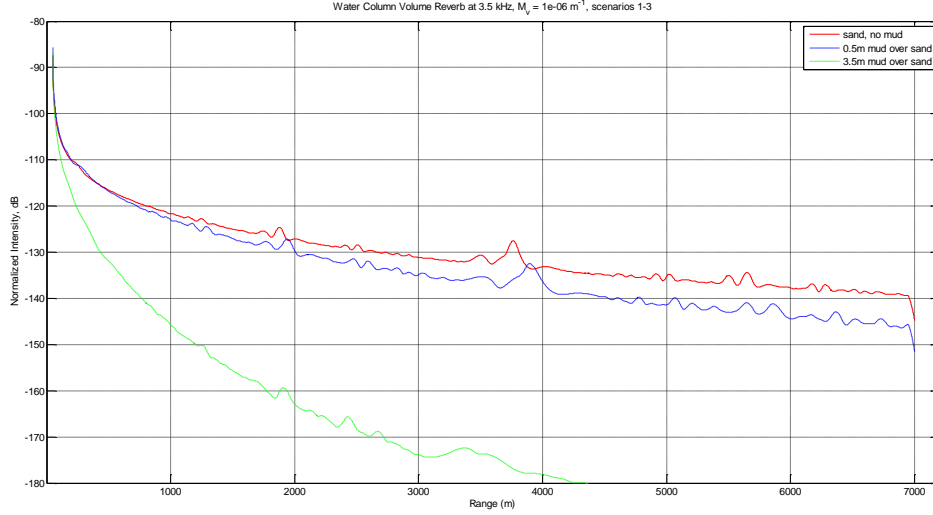


Figure 7: Normalized intensity of reverberation, $RL-SL$, at 3.5 kHz, caused by volume scattering in water column with $M_V = 10^{-6} \text{ m}^{-1}$. Shallow water (20m) with three different types of bottom: sand half-space (red), 0.5m of mud over sand (blue), and 3.5m of mud over sand half-space (green). The source is 1 m above bottom.

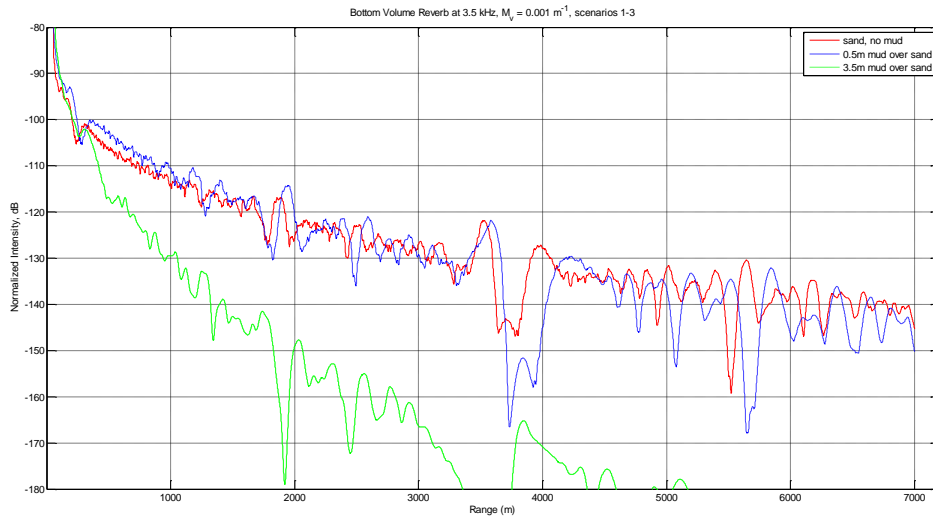


Figure 8: Normalized intensity of reverberation, $RL-SL$, at 3.5 kHz, caused by sediment volume scattering with $M_V = 10^{-3} \text{ m}^{-1}$. Shallow water (20m) with three different types of bottom: sand (red), 0.5m of mud over sand (blue), and 3.5m of mud over sand (green). The source is 1 m above bottom.

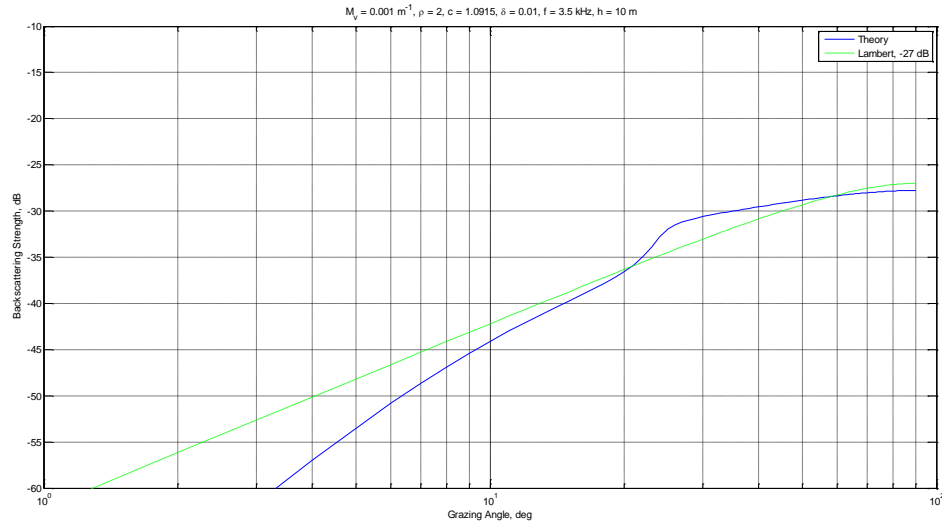


Figure 9: Angular dependence of Bottom Scattering Strength (Normalized Short-Range Reverberation Intensity), RL-SL, for 3.5 kHz, sand bottom with the same (typical) scattering strength per unit sediment volume used in Figure 8, $M_V = 10^{-3} \text{ m}^{-1}$. Also shown (in green) is Lambert Law at -27 dB.

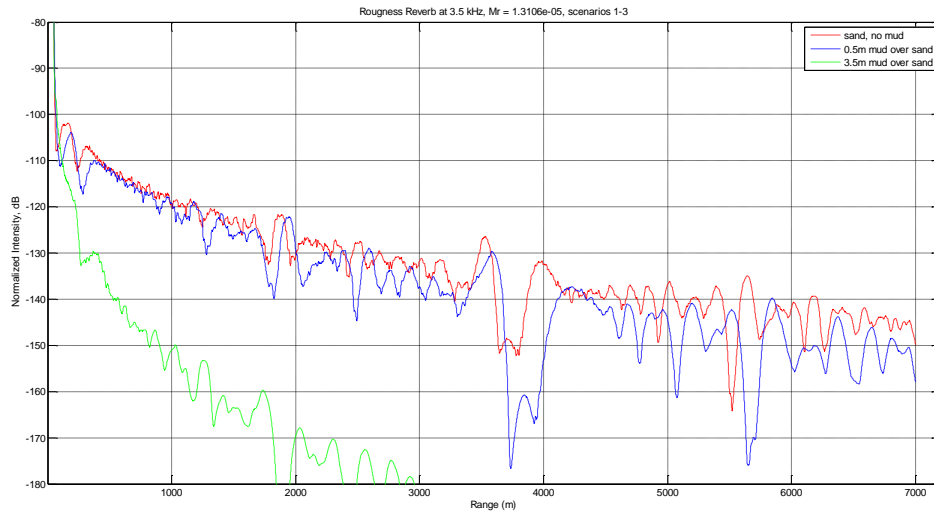


Figure 10: Normalized intensity of reverberation, RL-SL, at 3.5 kHz, caused by bottom interface roughness, with $M_R = 1.3 \times 10^{-5}$. Shallow water (20m depth) with three different types of bottom: sand half-space (red), 0.5m of mud over sand (blue), and 3.5m of mud over sand half-space (green). The source is 1 m above bottom.

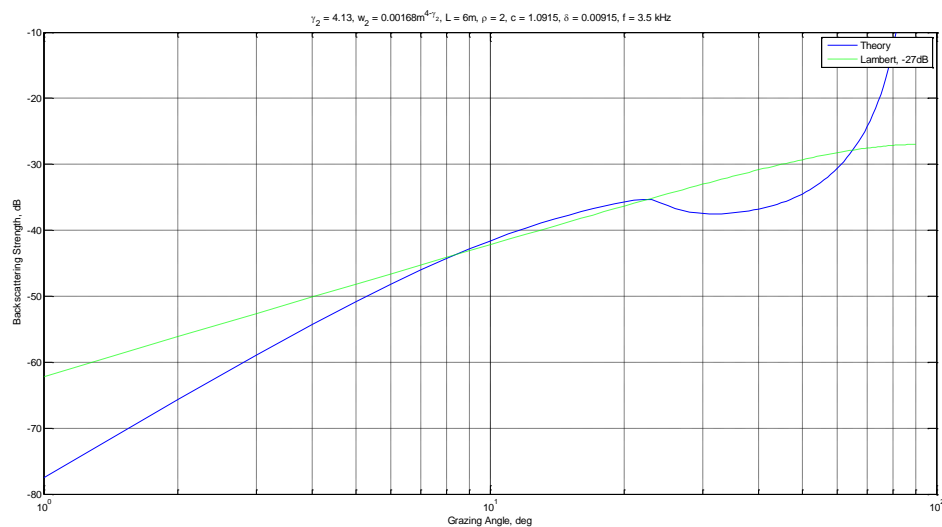


Figure 11: Angular dependence of Bottom Scattering Strength (Normalized Short-Range Reverberation Intensity), RL-SL, at 3.5 kHz, sand bottom with a roughness scattering coefficient used in Figure 10, $M_R = 1.3 \times 10^{-5}$. Also shown (in green) is Lambert Law at -27 dB.

Application to TREX Data Analysis

The described in previous sections approach was refined and specified for TREX13 environment, geometry and other input parameters.

TREX environmental scenario and input parameters

As in previous examples, Green's function magnitude was calculated using a PE code, Tang (2014), but for a slightly different configuration and parameters corresponding to a Pekeris waveguide with sand bottom. Parameters for these calculations were chosen as most relevant to TREX13 conditions and available by now acoustic data, Tang (2015):

Central frequency – 3450 kHz;
Frequency bandwidth – 100 Hz;
Water depth – 19 m;
Source depth – 17.8 m;
Receiver (FORA) depth – 16.9 m;
Azimuthal beamwidth – 2.6 deg;
Source Level (SL) – 200 dB.

Calculations for the penetration field within the sand bottom were made for depths down to 2 m below bottom surface.

Acoustic parameters for water and sand are taken as follows:

	water	sand
Sound speed, (m/s)	1525	1630
Density (g/cm ³)	1	2.0
Attenuation (loss tangent), δ	0	0.01

TREX-specified algorithm for numerical simulations

The algorithm for calculations of reverberation, specified for the TREX scenario, also includes three steps, but with some distinctions from those described earlier. The first step, in contrast with previous examples, considers a more general case of bistatic geometry, i.e., according to Eq.(2), now the propagator is to be calculated considering different locations of the source and receiver. For this, the Green's function magnitude was pre-calculated for refined input parameters, given above. Again, it is important that the algorithm requires calculations only at one (central) frequency.

The second step of the algorithm was also refined taking into account the fixed-fixed source-receiver geometry of the TREX13 reverberation measurements and a broad-band frequency filter (100 Hz) chosen for the TREX13 reverberation data processing, Yang et al (2014). This case requires a technique for smoothing strong oscillations of the propagation intensity, which differs from that used in previous examples. It takes into account a range-dependent scales in interferential structure of the broad-band field in shallow water waveguides. In this case the range span for the smoothing is range dependent as well, and, in contrast with previous numerical examples, is defined as follows: $L = r\Delta f / f$. More details on this technique are described in Ivakin (2015b).

The third step of the algorithm provides extremely fast evaluation of TREX13 reverberation using range-depth dependence of single-frequency Green's function intensity pre-calculated for this environment and smoothed in correspondence with the above-described technique. The integration, according to Eq.(2), was performed over a horizontal cross-section of the scattering (ensonified) volume, with a span of ranges, Δr , and a span of azimuthal angles, $\Delta\varphi$, defined in the same way as in previous examples, i.e. $\Delta\varphi = 2.6^\circ$ and $\Delta r = 8\text{ m}$.

TREX Reverberation: tentative model/data comparison

For further analysis, reverberation caused by bottom volume heterogeneity, as well as volume reverberation from scatterers in water column, was calculated and a tentative model/data comparison for a sub-set of TREX reverberation data was performed. The data sub-set used here is limited to a published example of TREX reverberation data, Hefner and Tang (2014), presented in some more detail in Yang et al (2014), see TREX13 Web-site. This example includes only one frequency sub-band ($f = 3400\div 3500$ Hz), one azimuthal look angle, $\varphi = 129^\circ$, and only two runs, Run #17 (23 April 2013) and Run #79 (9 May 2013). The data are shown in Figure 12, where a normalized reverberation intensity, Reverberation Level (RL) corrected for Source Level (SL), i.e. $RL - SL$, is presented.

It is seen that the data shown in Figure 12 are somewhat different. This difference might represent an effect of environmental changes resulted from a storm occurred in the area around 1 May 2013. The following calculations and results of model-data comparisons, presented in Figures 13-19, give a tentative interpretation of this effect.

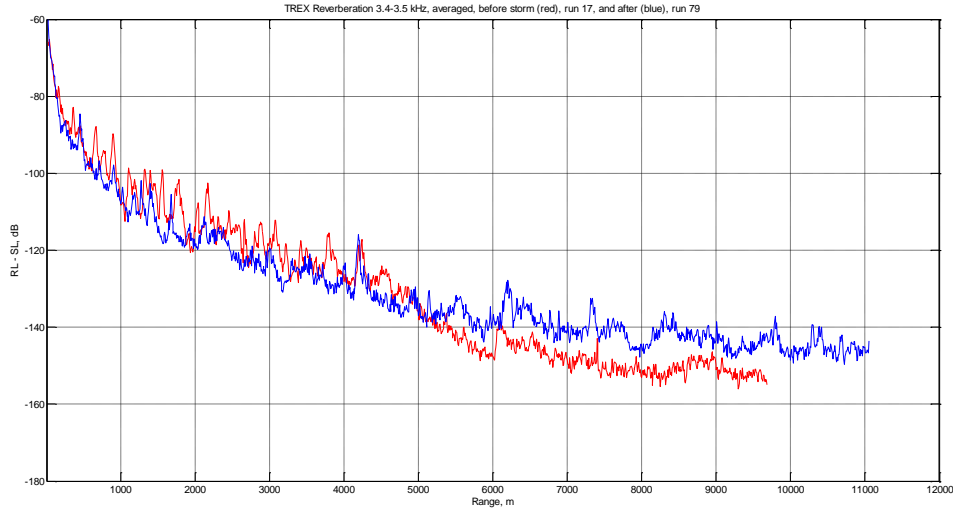


Figure 12: *TREX13 normalized reverberation intensity, $RL-SL$, at 3.4-3.5 kHz, for two runs, before (red) and after (blue) storm.*

The model suggests that reverberation is caused by volume scattering either in water column or in the sediment. The contribution of bottom roughness can be considered in a way similar to shown in previous numerical examples and is not included in this report. A separate paper, Ivakin (2015c), which discusses this case in detail, is in preparation. Typical values for the only free parameter of the model, M_v , were used to provide a reasonably good fit between model results and the data.

First, reverberation from water column was considered. The results for before storm model-data comparison is shown in Figure 13. Generally, it shows that scattering in water column can provide a good fit to data at long ranges (more than 6 km). However, to fit the data level, volume scattering strength in water column should be somewhat higher than typical. Unfortunately, there is no any available TREX ground truth data for scattering in water column at these ranges.

Next case to consider is volume scattering in the sediment. As an initial step on this way, the bottom volume scattering coefficient was assumed to be range-independent. The results are shown in Figure 14. Then a possibility of a smooth range-dependence of $M_v(r)$ was considered to reduce the model-data difference. This resulted in estimated/inferred range-dependence $M_v(r)$ shown in Figure 15. Corresponding corrected reverberation model-data comparison is shown in Figure 16.

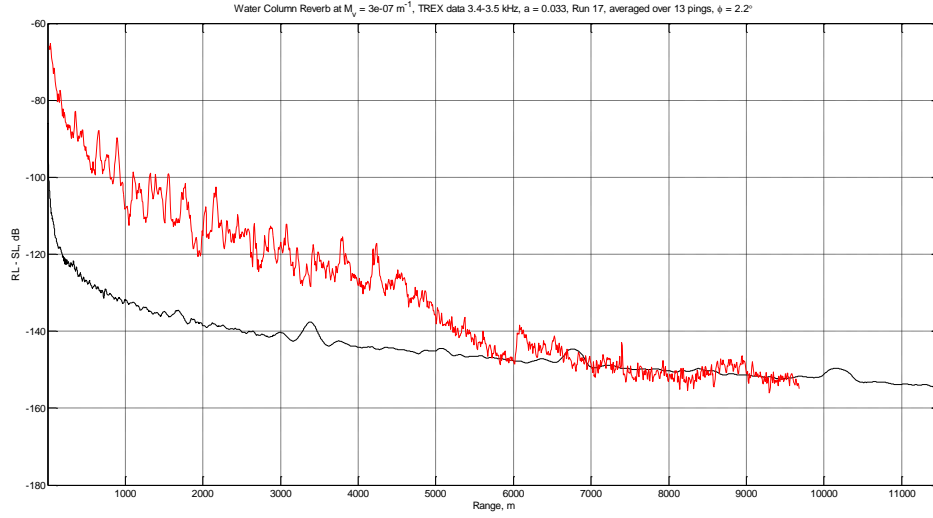


Figure 13: *TREX13 normalized reverberation intensity, $RL-SL$, at 3.4-3.5 kHz, measured before storm (red), compared to model result calculated for scattering in water column with volume scattering coefficient $M_v = 3 \times 10^{-7} \text{ m}^{-1}$.*

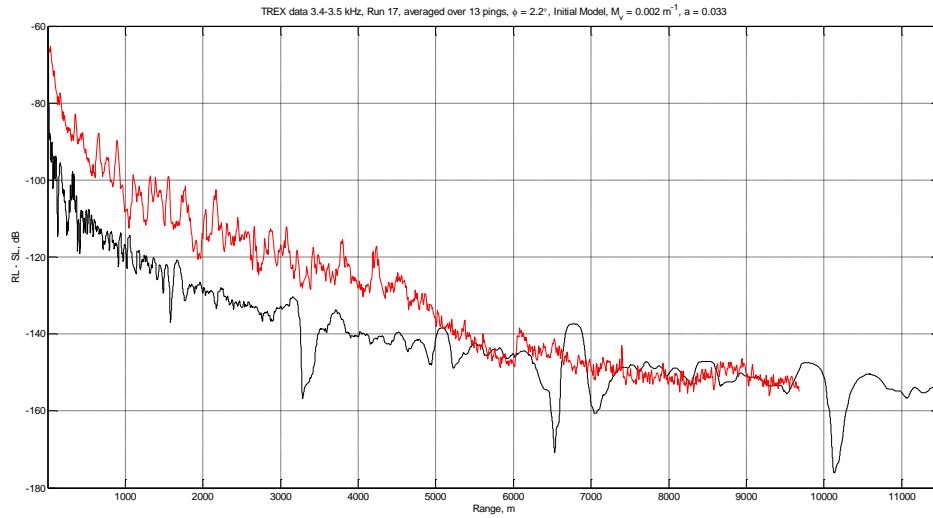


Figure 14: *TREX13 normalized reverberation intensity, $RL-SL$, at 3.4-3.5 kHz, measured before storm (red), compared to that calculated for volume scattering in sand bottom (range-independent) with $M_v = 0.002 \text{ m}^{-1}$.*

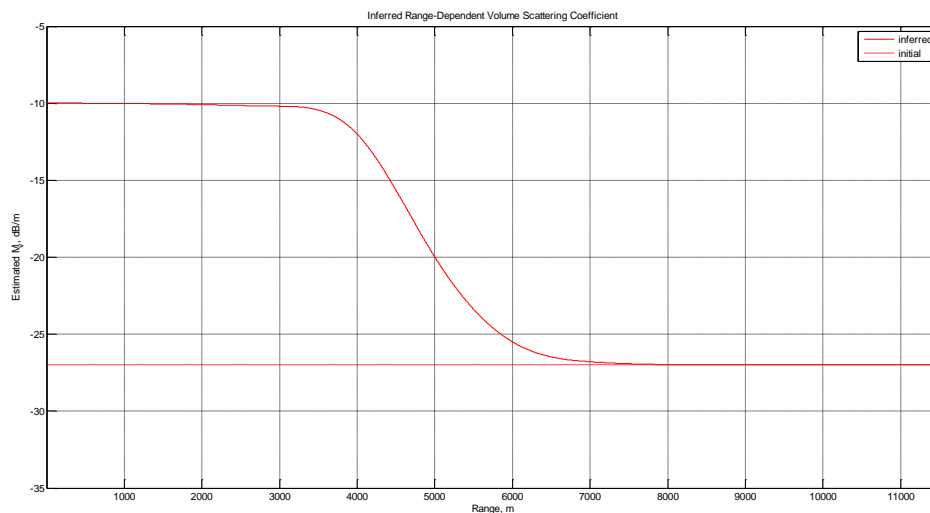


Figure 15: A correction made for range-dependence of local volume scattering strength in TREX13 sand bottom before storm.

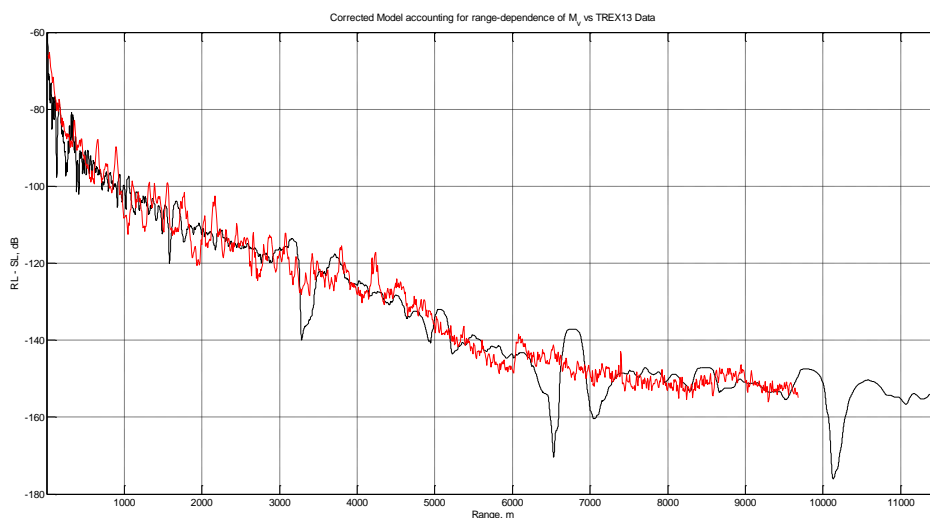


Figure 16: TREX13 normalized reverberation intensity, $RL-SL$, at 3.4-3.5 kHz, measured before storm, compared to that calculated for volume scattering in sand bottom with corrected range-dependent volume scattering strength shown in Figure 15.

A similar procedure was performed for analysis of after storm reverberation. The results for initial model data comparison (for range independent M_V) are shown in Figure 17. It shows only a slight discrepancy at small ranges, which was easily compensated by a correspondingly slight range-dependence of sediment volume scattering coefficient shown in Figure 18. Then model-data comparison, shown in Figure 19, becomes reasonably good.

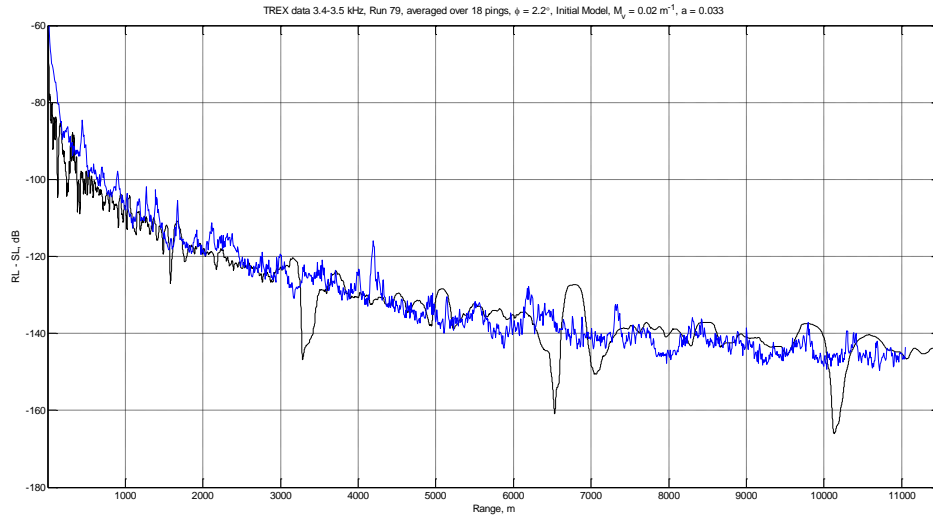


Figure 17: *TREX13 normalized reverberation intensity, $RL-SL$, at 3.4-3.5 kHz, measured after storm (blue), compared to that calculated for volume scattering from sand bottom (range-independent),) with $M_V = 0.02 \text{ m}^{-1}$.*

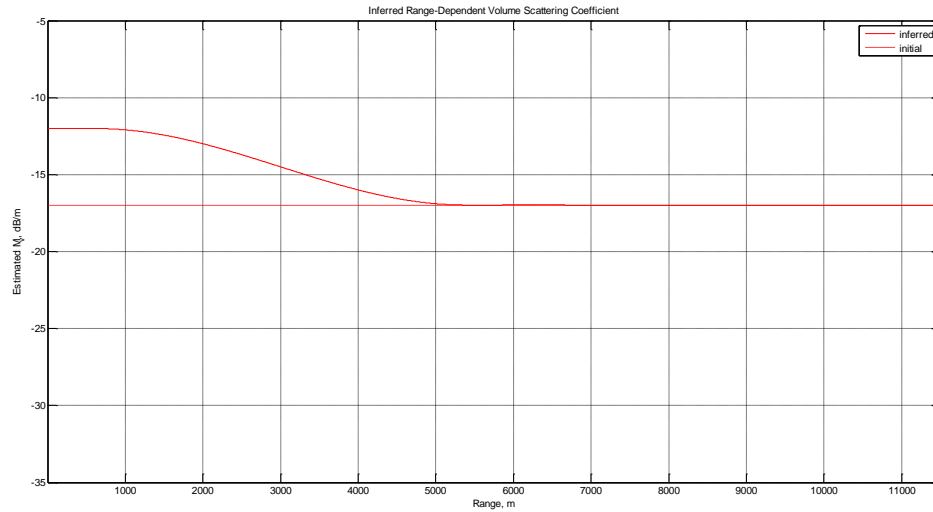


Figure 18: *A correction made for range-dependence of local volume scattering strength in TREX13 sand bottom after storm.*

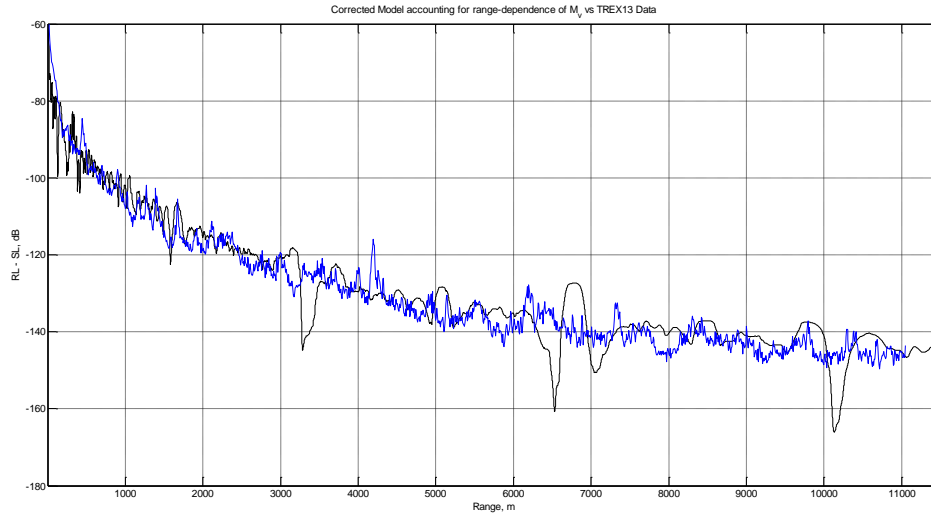


Figure 19: *TREX13* normalized reverberation intensity, $RL-SL$, at 3.4-3.5 kHz, measured after storm (blue), compared to that calculated for volume scattering in sand bottom with corrected range-dependent local volume scattering strength shown in Figure 18.

Comparison of Figure 15 and Figure 18 may help to understand nature of the environmental impact of the storm in the area that significantly affected reverberation. Such impact could result from dynamical mixing of the sediment and consequent smoothing of the range dependence of the strength of surficial sediment heterogeneity along the observation path. Importantly, same conclusions would be reached if the range-dependence of local roughness scattering is considered using Eq.(5). In more detail, this will be described in a separate paper, Ivakin (2015c).

In addition, to ensure that the two runs, #17 and #79, see Figure 12, used here for model-data comparisons, are indeed representative of a more comprehensive data set, reverberation data from five more runs which recently became available, Tang (2015), were considered. They are added to previous two runs and shown in Figure 20 and confirm the general effect of the storm seen in Figure 12. This suggests that above-shown model-data comparison, as well as conclusions based on this comparison, will stay reasonable for a more comprehensive TREX13 data analysis.

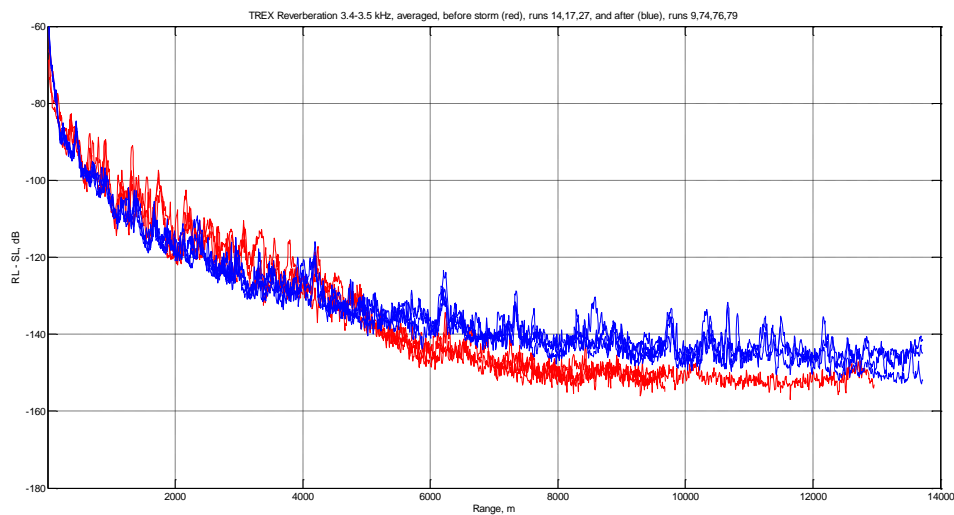


Figure 20: *TREX13 normalized reverberation intensity, $RL-SL$, at 3.4-3.5 kHz, measured before (red) and after (blue) storm. Five more runs added to Figure 12, showing that previously used two runs are typical.*

Concluding Remarks and Implications for Future Research

This report has presented a modeling approach that allows fast estimations of volume reverberation in complex shallow water environments. Computer simulations with environmental inputs typical for shallow water have been conducted. Capabilities of this modelling approach and developed codes to provide a reasonable interpretation of TREX2013 reverberation data are demonstrated and discussed.

Based on these results, future research may address several science issues, which appeared to be important for better understanding TREX2013 data:

- To account for sediment lateral variability, mud patches/strips, using more advanced PE codes for Green's function (here only the case of Pekeris waveguide was considered);
- To analyze environmental and acoustic data to get inputs describing variability of shell content, sub-bottom interface spectra, heterogeneity in mud;
- To extensively support TREX13 data analysis, including data obtained for whole mid-frequency range, 1-10 kHz, and all other azimuthal angles (here only a subset of data at 3.4-3.5 kHz and one azimuthal direction was considered).

Several journal papers are planned to be prepared for submission in near future, based on results and ideas of this research, briefly discussed and presented in this report. One of them, Ivakin (2015b), soon to be submitted, includes results shown in Figures 12-19 of this report and use data currently available on TREX2013 web-site (Figure 12). More papers are in preparation, tentatively shown as Ivakin (2015c, 2015d, 2015e).

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- A.N. Ivakin (2013), "Volume scattering and reverberation in shallow water: A simplified modeling approach", *J. Acoust. Soc. Amer.*, 134 (5, Pt.2), p.4208, 2013.
- B.T. Hefner, D.R. Jackson, A.N. Ivakin, and G. Wendelboe (1913), "Physics-based inversion of multibeam sonar data for seafloor characterization", *J. Acoust. Soc. Amer.*, **134**(4), Pt.2, p.4240.
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- A.N. Ivakin (2015), "A modeling approach to volume scattering and reverberation in shallow water", in, *Proc. 3rd Underwater Acoustics Conference and Exhibition*, J.S. Papadakis and L. Bjorno (Eds), Crete, Greece, 21-26 June 2015, p.53.

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- Ivakin A.N. (2015d), “Scattering from discrete heterogeneity of marine sediments”, in preparation.
- Ivakin A.N. (2015e), “A sediment characterization technique based on measurements of long-range reverberation in shallow water”, in preparation.
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1. REPORT DATE (DD-MM-YYYY) 31-08-2015		2. REPORT TYPE Final Report		3. DATES COVERED (From - To) 1 December 2012 - 31 May 2015	
4. TITLE AND SUBTITLE Modeling of Mid-Frequency Reverberation in Very Shallow Water: A Greens FUnction Approach and Application to TREX2013 Data Analysis				5a. CONTRACT NUMBER	
				5b. GRANT NUMBER N00014-10-1-0332	
				5c. PROGRAM ELEMENT NUMBER	
				5d. PROJECT NUMBER	
6. AUTHOR(S) Anatoliy N. Ivakin Applied Physics Laboratory University of Washington				5e. TASK NUMBER	
				5f. WORK UNIT NUMBER	
7. PERFORMING ORGANIZATION NAME(S) AND ADDRESS(ES) Applied Physics Laboratory University of Washington 1014 NE 40th St. Seattle, WA 98105				8. PERFORMING ORGANIZATION REPORT NUMBER	
9. SPONSORING/MONITORING AGENCY NAME(S) AND ADDRESS(ES) Office of Naval Research, Code 322 875 North Randolph Street Arlington, VA 22203-5320				10. SPONSOR/MONITOR'S ACRONYM(S) ONR	
				11. SPONSOR/MONITOR'S REPORT NUMBER(S)	
12. DISTRIBUTION/AVAILABILITY STATEMENT Distribution Statement A: Approved for public release, distribution unlimited.					
13. SUPPLEMENTARY NOTES					
14. ABSTRACT The long-term goals of this research are to better understand and accurately model low- to mid-frequency reverberation in shallow water environment. Specific goals are to develop a model of reverberation for conditions (1 – 10 kHz, ~20 m water depth, ~ 10 km range) corresponding to the ONR Target and Reverberation Experiment performed on the spring 2013 (TREX2013), develop a code and conduct computer simulations with environmental inputs typical for the chosen location, and apply this model to analysis of available TREX2013 data. This report presents a Green's function modeling approach that allows fast estimations of volume reverberation in complex shallow water environments. A simplified first-order version of the approach is considered to show how far-field scattering solutions obtained for free space can be incorporated into reverberation in complicated bounded, range-dependent, and stratified environments. A higher-order modification of this approach is considered as well, using a MFSB (Multiple Forward Single Backscatter) approximation. Application to analysis of TREX2013 reverberation and tentative model-data comparisons are presented.					
15. SUBJECT TERMS Underwater acoustic propagation and reverberation, Green's function in a shallow water waveguide, volume and roughness scattering, data inversion, TREX13 data analysis and model-data comparisons.					
16. SECURITY CLASSIFICATION OF:			17. LIMITATION OF ABSTRACT U	18. NUMBER OF PAGES 33	19a. NAME OF RESPONSIBLE PERSON Anatoliy N. Ivakin
a. REPORT U	b. ABSTRACT U	c. THIS PAGE U			19b. TELEPHONE NUMBER (Include area code) 206-616-4808